

# The multiscale entanglement renormalization ansatz (MERA) (80)

MPS: very well suited for gapped 1D systems.

→ area law

→ exp. decay of correlations (length scale!)

What if system is gapless (i.e. at phase transition)?

→ diverging correlation length

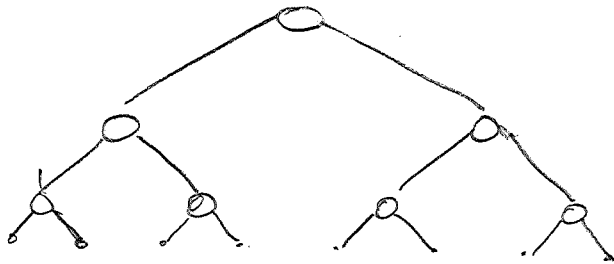
→ system is scale invariant  $\leftrightarrow$  no length scale

→ correlations: algebraic (corr  $\sim e^{-\alpha}$ )

→ logarithmic scaling of block entropy:  $S(S_L) \sim \log L$ .

Ansatz for scale-invariant systems?

First idea: tree tensor networks:

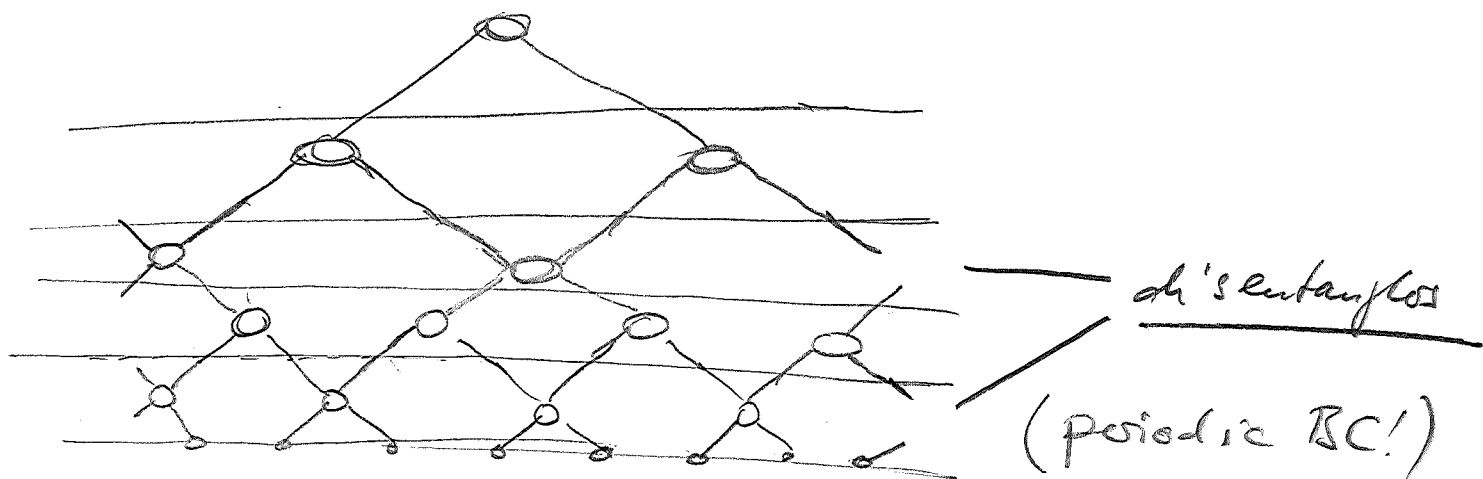


⊕ Has scale invariance.

⊖ Entanglement: still area law scaling

⊖ Entanglement detn. closely requires dep. on position hard to establish.

→ Include "disentanglers" which disentangle regions: (81)



We will choose the disentanglers to be unitary  
and the other tensors to be isometries.

"Multiscale entanglement renormalization ansatz" (MERA)

(G. Vidal, quant-ph/0610099)

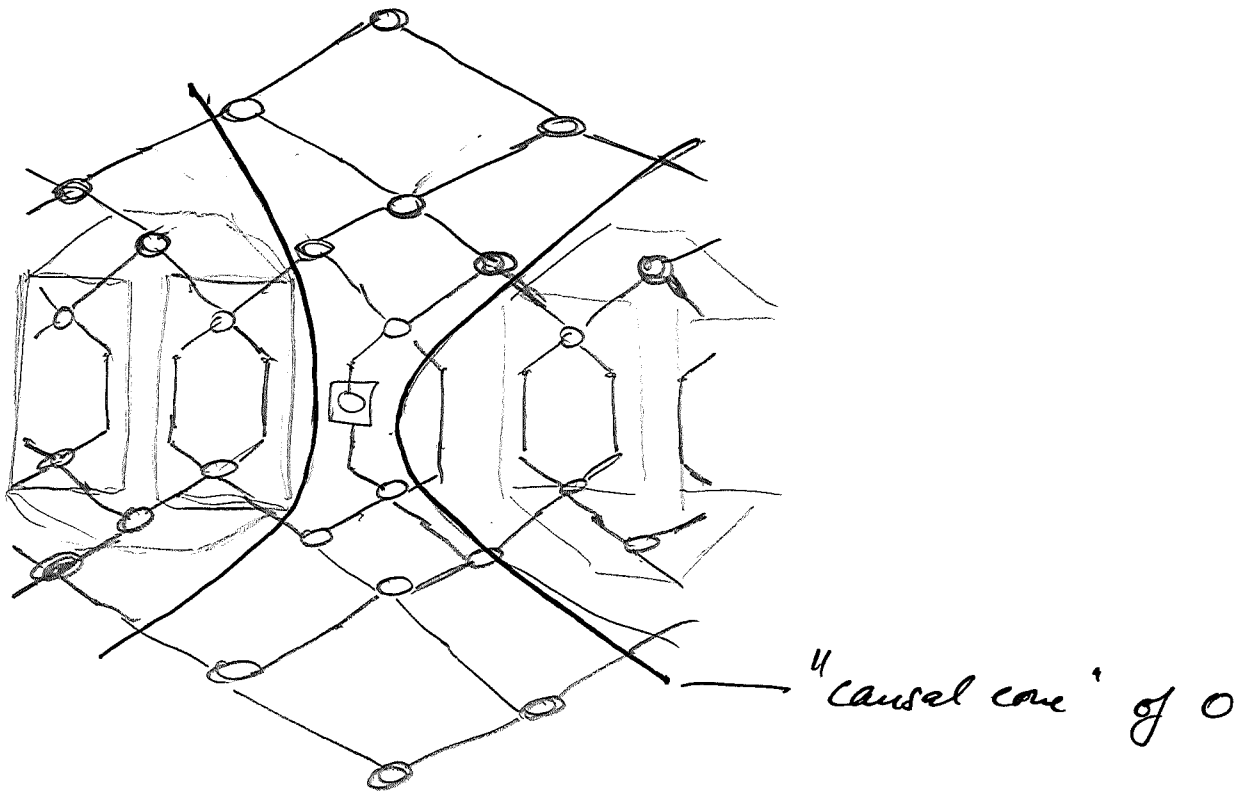
We use some maximal dimension  $\chi$  for the tensor indices.

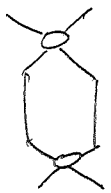

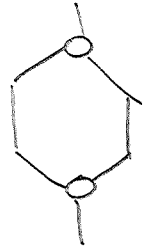

→ Efficient parametrization (in  $\chi$  and  $N$ ) of a class of 1D states (w/ built-in scale invariance) — scales like

$$O(\chi^4 \cdot N)$$

By choosing identical tensors in all layers, we obtain a scale-invariant state!

Can we compute exp. values efficiently?



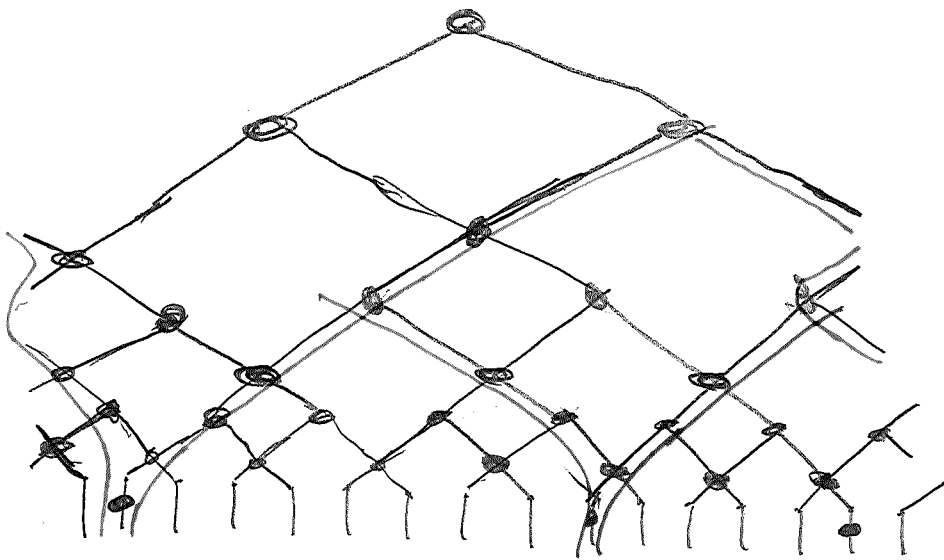
use  =  and  = 

to cancel most of the network!

$\Rightarrow$  "causal cone" of operator  $O$  has always fixed width (in terms of tensors! - in "real space", the width grows exponentially!), and depth  $\log N$ .

$\Rightarrow$  Can be contracted efficiently! (cost:  $\log N \cdot \text{poly}(X)$ )

# Correlation functions:



etc.

two causal cones of fixed width, which eventually fuse to one causal cone

⇒ correl. functions can be computed efficiently!

What scaling should we expect for correl. functions?

- the two causal cones intersect after  $\log l$  layers.

↑ distance

- If tensors are the same (scale inv.), each layer reduces correlations by some factor  $\lambda$ .

$$\Rightarrow \text{Corr} \sim \lambda^{2 \log l} = l^{2 \log \lambda}$$

⇒ algebraic decay of correlations expected!

What entanglement scaling do we expect?

84

Entanglement  $\equiv$  # bonds cut = # layers (depth)

$\Rightarrow$  Ent. should scale like  $\log L!$

Simulation w/ MERA:

Variationally optimize energy as function of tensors.

Either: all tensors identical (per layer or all),  
or different tensors everywhere.

Note: Although energy is bilinear in each tensor, this  
cannot be mapped to eigenvalue problem due to  
unitarity constraint.

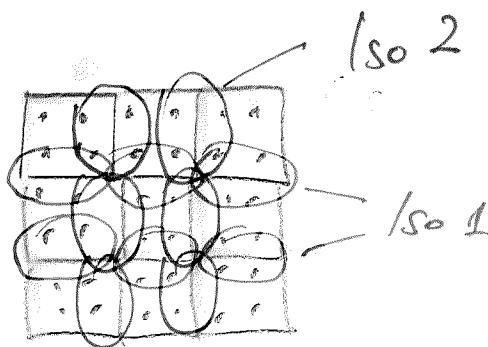
Higher dimensions:

Can be generalized to 2D etc.

Idea always: Partition system  
into blocks, disentangle b/w

blocks, and reduce #

blocks w/ isometries:



Note: Contraction still exact - no approx. needed (cf. PEPS)

Unitary tensor networks:

Basic idea for efficient contraction in TERA:

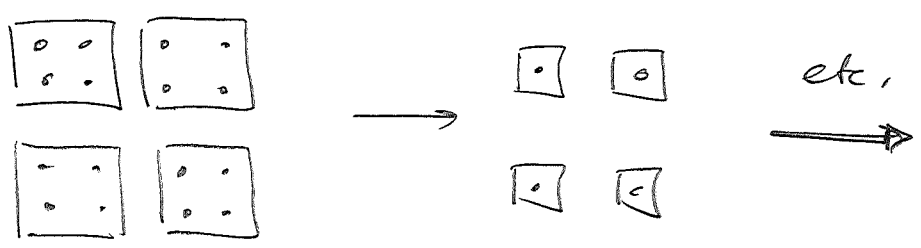
Unitarity of tensors

This can be used to construct a large class of contractible tensor networks where the tensors are unitary!

Relation to renormalization (RG):

Real space renormalization (Kadanoff, Wilson, ...):

Block sites in a system and discard "irrelevant local" degrees of freedom.



E.g. Ising model: Majority vote!

(This also gives a trace of Hamiltonian & temperature for the "effective Ham." on renormalized deg. of freedom)

This procedure isolates the long-range features of the system & removes short-range fluctuations.

E.g. Ising model, low temperature: after a few RG steps, all effective spins point in the same direction  $\rightarrow$  RG decreases temperature.

For high temp., the pattern will become more random by RG  $\rightarrow$  T increases.

Critical states/systems: invariant under RG, since no length scale.

Problem w RG for quantum states: IS/1

\* To discard irrelevant DoF, we want to rotate them, i.e., they should factor.

\* Irrelevant local DoF will however also be entangled across blocks

$\Rightarrow$  need for disentanglers!