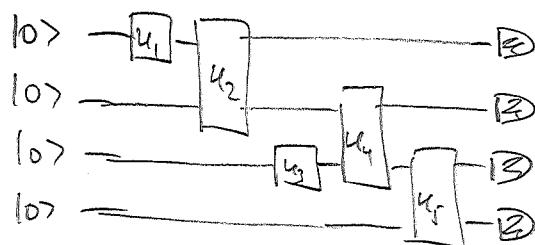


Measurement based quantum computation

(Review: quant-ph/0508124)

Standard model of quantum computation: Circuit model:



* N -qubit "quantum register" $(\mathbb{C}^2)^{\otimes N}$

1. initialize to $|0\rangle^{\otimes N}$

2. apply one- and two-qubit "gates" (= unitaries) from

"universal gate set", e.g. $R_x(\phi) = e^{-i\phi X}$, $R_z(\phi) = e^{-i\phi Z}$,

$$CP = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad (\text{"Controlled-phase"})$$

3. measure (a subset of) qubits in 2-basis $\{|0\rangle, |1\rangle\}$

→ output of computation!

Requires to keep register coherent and ability to apply two-qubit unitaries (and interactions are typ. complicated)

→ alternative models: adiabatic QC, topological QC, dissipative QC, measurement based QC, ...

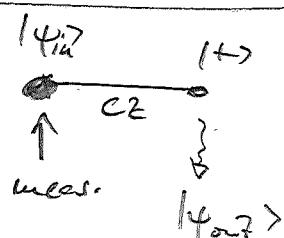
Researched based q. computation (RBQC):

1. Prepare special 2D state ("cluster state"): ground state of gapped local ham!
2. Perform a sequence of (adaptive) one-qubit meas.
3. Output of QC function of meas. outcome
 → No interactions needed after cluster state is prepared!

(Note: Also known as "one-way quantum computing")

Elementary "gadget":

- Consider $|+\rangle_1 |+\rangle_2$: ($|+\rangle = \alpha|0\rangle + \beta|1\rangle$)
- Apply a C2
- measure 1 in basis $\left(e^{+i\phi/2}|0\rangle \pm e^{-i\phi/2}|1\rangle \right)/\sqrt{2}$

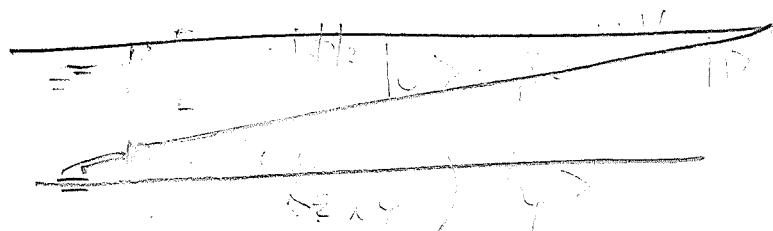


Outcome:

$$|\psi_{in}\rangle |+\rangle_2 = \alpha|0\rangle |+\rangle + \beta|1\rangle |+\rangle$$

$$\xrightarrow{\text{C2}} \alpha|0\rangle |+\rangle + \beta|1\rangle |-\rangle$$

$$\xrightarrow{\text{meas.}} |\psi_{out}\rangle = \alpha e^{-i\phi/2} |+\rangle \pm \beta e^{+i\phi/2} |-\rangle$$



with m the measurement outcome,

and $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$ the Hadamard gate, we find...

Note: H changes Schw. X and Z basis:

$$H|+\rangle = |0\rangle, H|0\rangle = |+\rangle$$

$$H|-\rangle = |1\rangle, H|1\rangle = |-\rangle$$

$$HXH = Z, HZH = X$$

$$H^2 = \mathbb{1}$$

$$\begin{aligned} |\psi_{out}\rangle &= \alpha e^{-i\phi/2} |+\rangle \pm \beta e^{+i\phi/2} |-\rangle \\ &= H [\alpha e^{-i\phi/2} |0\rangle \pm \beta e^{+i\phi/2} |1\rangle] \\ &= H \overset{\text{“}}{Z} R_z(\phi) |\psi_{in}\rangle \\ &= \overset{\text{“}}{X} H R_z(\phi) |\psi_{in}\rangle \end{aligned}$$

\Rightarrow Protocol implement 1-qubit-gate $HR_z(\phi)$ up to a Pauli error!

Concatenate two steps:

Step 1: Angle $\phi_1 \rightarrow$ Outcome m_1 ,

Step 2: Angle $(-1)^{m_1} \phi_2 \rightarrow$ Outcome m_2

$$\begin{aligned} \Rightarrow |\psi_{out,2}\rangle &= \overset{\text{“}}{X}^{m_2} \underbrace{HR_z((-1)^{m_1} \phi_2)}_{= \overset{\text{“}}{X}^{m_1} R_z(\phi_2)} \overset{\text{“}}{X}^{m_1} HR_z(\phi_1) |\psi_{in}\rangle \\ &= \overset{\text{“}}{Z}^{m_1} H \end{aligned}$$

$$= X^{u_2} Z^{u_1} \underbrace{H R_z(\phi_2) H}_{= R_x(\phi_2)} R_z(\phi_1) | \psi_n \rangle$$

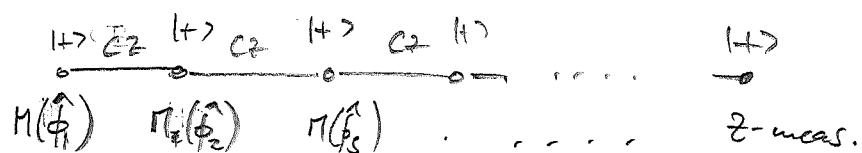
$$= X^{u_2} Z^{u_1} R_x(\phi_2) R_z(\phi_1) | \psi_n \rangle.$$

Can be iterated:

$$\begin{aligned} & X^{u_4} Z^{u_3} R_x(\tilde{\phi}_4) R_z(\tilde{\phi}_3) X^{u_2} Z^{u_1} R_x(\phi_2) R_z(\phi_1) | \psi_n \rangle \\ &= X^{u_4 \oplus u_2} Z^{u_3 \oplus u_1} \underbrace{R_x((-1)^{u_1} \tilde{\phi}_4)}_{\phi_4} \underbrace{R_z((-1)^{u_2} \tilde{\phi}_3)}_{\phi_3} R_x(\phi_2) R_z(\phi_1) | \psi_n \rangle \end{aligned}$$

Read-out: We can simply measure in the Z basis (and correct for the X error!).

Protocol for 1-qubit computation; starting in state $|+\rangle$:



But: We can as well first do all CZ and then measure.

Moreover, the CZ commutes - order irrelevant.

→ "cluster state"

This is the unique ground state of

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$$H = -\sum h_i, \quad h_i = \sum_{j=i} z_{j-1} \otimes x_j \otimes z_{j+1}$$

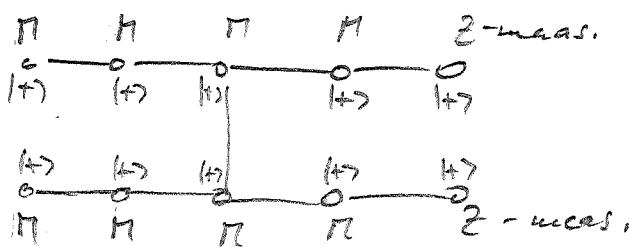
(Proof: Homework)

→ Computation based on G.S. & 1-qubit-meas.

How can we go beyond 1 qubit?

$$\begin{array}{c} \text{two-qubit } \{ \circ \\ \text{input } |\psi_{in}\rangle \{ . \end{array} \xrightarrow{\text{CZ}} \begin{array}{c} \bullet \\ \text{CZ} \end{array} \longrightarrow |\psi_{out}\rangle = \text{CZ } |\psi_{in}\rangle \\ (\text{without measurement!}) \end{array}$$

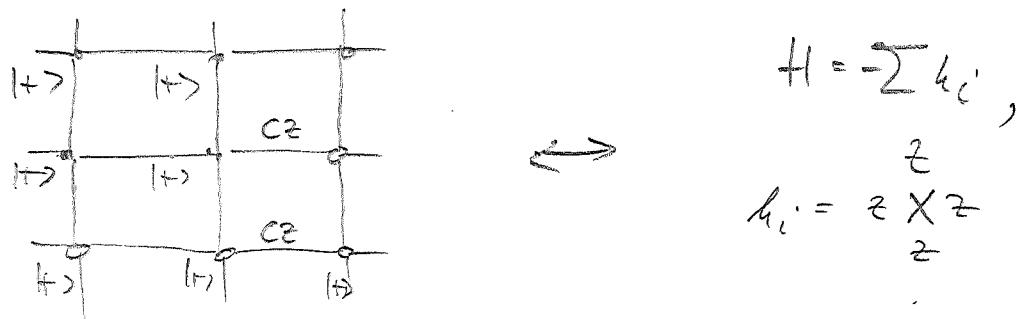
→ two-qubit operations in a circuit via



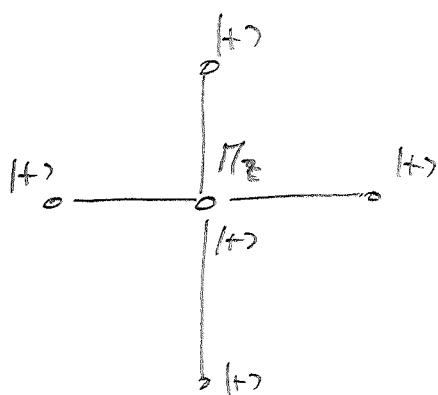
⇒ Can again be based on cluster state on the right underlying graph! (This is again G.S. of L.H. $H = -\sum h_i$ with $h_i = \bigotimes_{j \in \text{neigh}(i)} z_j \rightarrow \text{homework!}$)

→ Any q.computation can be implemented by first preparing a cluster state & then doing 1-qubit-meas!

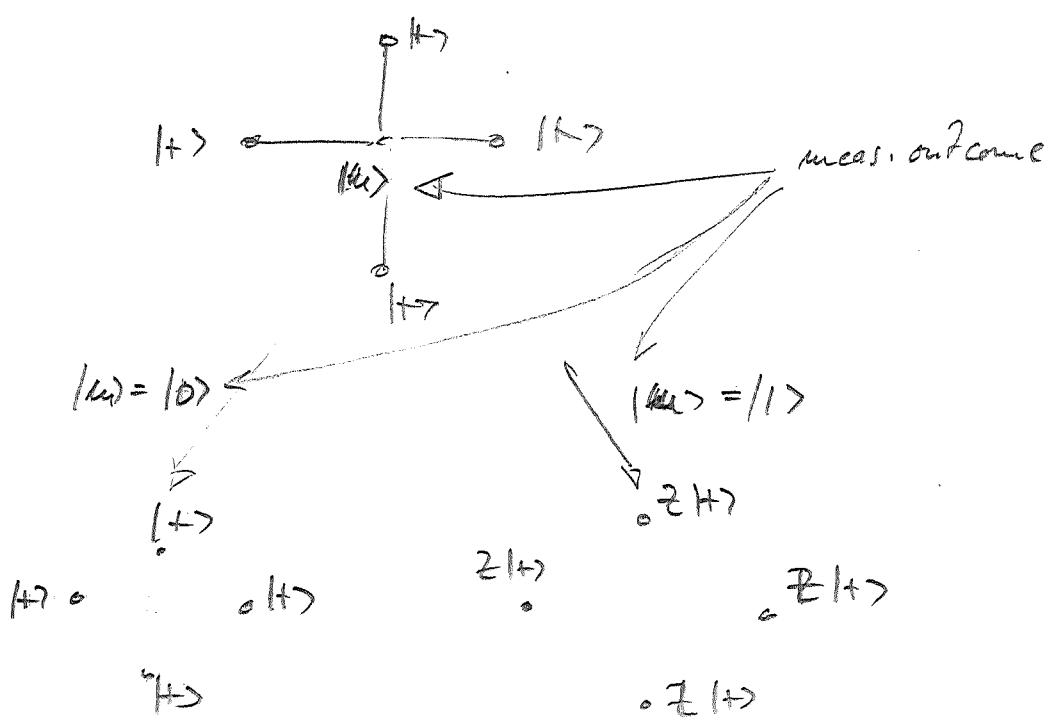
Can we also do this on e.g. a square lattice:



→ 2-meas. allow to erase sinks from a cluster state:



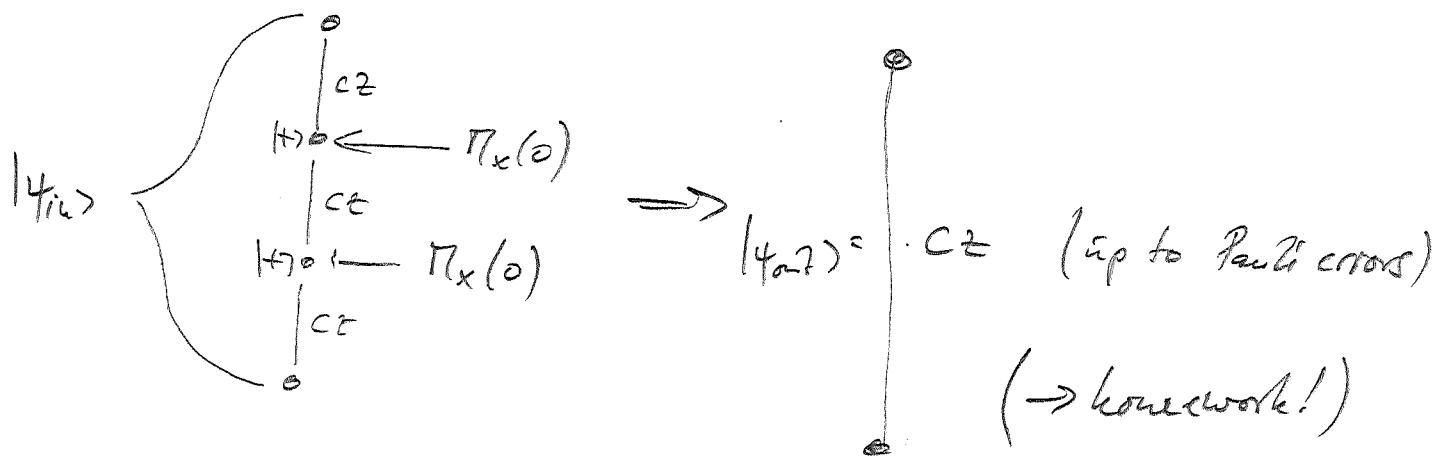
Meas. commutes w/ CZ → equiv. to



→ qubit removed of 2-errors or neighbors

→ 2-err. correct w/ CZ → we can use this to "etch" a circuit² into a regular (e.g. 2D square lattice) clock state! (only need to adapt meas. bases to 2 errors)

But: We need a different way to do 2-qubit gates (on the square lattice, at least):



Full protocol:

- Start from clock state
- "etch" circuit via 2 meas.
- perform sequence of XY-plane-meas. to implement circuit (adaptive SGS!)
- measure output on 2 basis & interpret according to prob. meas. outcomes ($\equiv X/2$ errors)

Remarks:

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- adaptive meas. only requires comp. of patterns → computationally very easy
- many meas. patterns can be implemented non-adaptively
- in part.: any Clifford circuit (generated by $\{H, S, CNOT\}$, with $S = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$) can be implemented by non-adaptive measurements
- non-adaptive measurement can be done in parallel - potential parallelization of Q.Comp.
- in certain cases, a logarithmic number of layers is enough.
(Note: class. side-provacy still requires poly time!)
- MBQC has a particularly nice interpretation in terms of teleportation and PEPS (\rightarrow homework)

A beautiful application of MBQC: "Blind q. computation" (Broadcast, Fiberman, Beaufit; arXiv: 0807.6454)

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Bob has full QC, Alice can only prepare single qubits. Can Bob perform a QC for Alice w/out know comp. & outcome?

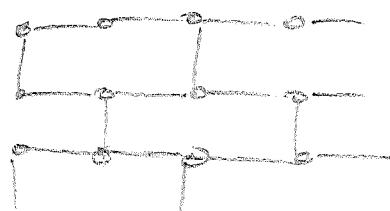
Idea: A prepares $|+\rangle$ states and sends them to Bob, who entangles them w/ C7 & performs the meas. A tells him.

Trick: A prepares $|+\phi\rangle = |0\rangle + e^{i\phi}|1\rangle$ instead w/ different random ϕ , for each qubit \rightarrow cluster has random Z rotations at each site.

Alice can adapt the meas. basis to the rotation, while for Bob, the meas. looks completely random. Bob reports outcomes & Alice adapts the meas. \rightarrow no info revealed!

The final Z meas. is also random to Bob since he does not know the X errors!

Bob could still learn sth. from the shape of the cluster \rightarrow use universal "bridewort state"



which can support any Q.C.!