

Interaction of quantized light w/ two-level system: Rabi oscillations

No detuning ("resonant"):

Cavity in Fock state:  $|q(t)\rangle = \cos\left(\frac{\omega_a t}{2}\right) |e, u\rangle + \sin\left(\frac{\omega_a t}{2}\right) |g, u+1\rangle$

with  $n$ -photon Rabi frequency  $\omega_a = \sqrt{n+1} \omega_0$  ( $\omega_0 = \frac{2dE_0}{\hbar}$ )

Cavity in arbitrary state  $\sum c_u |u\rangle$ , atom in  $|e\rangle$ :

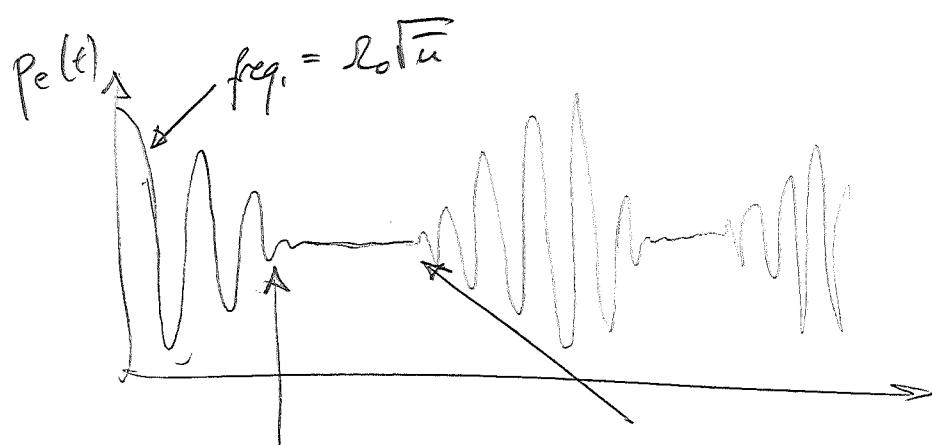
$$|q(t)\rangle = \sum_u c_u \left[ \cos \frac{\omega_0 \sqrt{n+1}}{2} t |e, u\rangle + \sin \frac{\omega_0 \sqrt{n+1}}{2} t |g, u+1\rangle \right]$$

Probability  $p_e(t)$  for atom in  $|e\rangle$ :

$$p_e(t) = \sum_{u>0}^{\infty} |c_u|^2 \cos^2\left(\frac{\omega_0 \sqrt{n+1}}{2} t\right) = \sum_{u>0}^{\infty} |c_u|^2 \frac{1 + \cos(\omega_0 \sqrt{n+1} t)}{2}$$

$\Rightarrow p_e(t)$  is sum of many incommensurate Rabi oscillations!

(Numerical) observation for coherent light  $|\alpha\rangle = \sum c_u |u\rangle$ :



collapse - could  
be explained  
classically  
(damping)

revival of Rabi oscillations  
- cannot be explained classically!

Explanation: For coherent light:

$$\bar{u} = |x|^2 ; \quad \Delta u = \sqrt{u} \quad (\text{Poisson distribution})$$

→  $\Delta u$  around  $\bar{u}$  with width  $\Delta u$ .

$$\cos(\Omega_0 \sqrt{u} t) = \cos\left(\Omega_0 \sqrt{u+1} \sqrt{1 + \frac{\delta u}{\bar{u}+1}}\right) \approx \cos\left(\underbrace{\Omega_0 \sqrt{\bar{u}+1}}_{=: \Omega} \left(1 + \frac{\delta u}{2\bar{u}+2}\right)\right)$$

$$= \cos \Omega \left(1 + \frac{\delta u}{2\bar{u}+2}\right) t \quad (\delta u = u - \bar{u})$$

→ initial oscillation w/ freq.  $\Omega = \Omega_0 \sqrt{\bar{u}+1}$

→ Oscillators get out of phase at rate  $\frac{\delta u}{2\bar{u}+2} t$

Collapse: After time  $\Omega \cdot \frac{\delta u}{2\bar{u}} t \sim \pi$ , oscillators are roughly

$u \pm \Delta u$  completely out of phase  $\rightarrow \Omega_0 \sqrt{\bar{u}} \frac{\pi}{2\bar{u}} t \sim \pi$

$\Rightarrow$  Collapse after  $t \sim \frac{2\bar{u}}{\Omega_0}$ : (independent of  $\bar{u}$ !)

→ approx.  $\sqrt{u}$  oscillations before collapse.

Revival: Oscillators back in phase if  $\Omega \cdot \frac{1}{2\bar{u}} t \sim 2\pi$ , i.e.

$$\Rightarrow \text{Revival after } t = \frac{4\pi}{\Omega_0} \sqrt{u}$$

Classical limit:  $\bar{u}$  very large,  $\Omega_0$  very small

$\rightarrow \Omega = \Omega_0 \sqrt{u} ; \quad \sqrt{u} \gg 1$  oscillations

Experimental realization: Rydberg atoms in micro-wave cavities (Haroche group; Nobel prize 2012)

### \* Two-level system:

Rubidium atom in circular Rydberg state  $|nC\rangle$ :

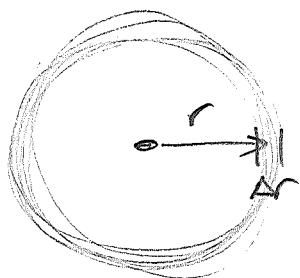
- 1 valence electron - Hydrogen-like

- valence electron lightly excited:

principal q. number  $n \approx 50$

max. angular momentum  $l = n - 1$  ( $\Rightarrow$  circular),  
magnetic q. numbers  $m = l$ .

Semi-classical orbit of electron:



circular orbit w/ radius  $r = a_0 n^2$ ,  
width of orbit  $\Delta r \sim 1/\sqrt{2n}$

Böhr radius,  $a_0$

States:  $|g\rangle = |50C\rangle$

$|e\rangle = |57C\rangle$

Energies:  $E_n = -\frac{R}{n^2} \Rightarrow \omega_{\text{eg}} = \frac{2R}{n^3}$

$$\Rightarrow \omega_{\text{eg}}/2\pi = 57 \text{ GHz} \quad (\lambda = 5.9 \text{ mm})$$

Matrix element of dipole transition  $\langle uC | \vec{r} | u+1, C \rangle$ :

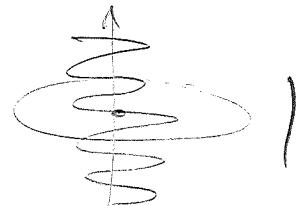
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Spatial overlap of  $|uC\rangle$  &  $|u+1, C\rangle$  almost 1 ( $\frac{r_{u+1}}{r_u} \sim \frac{2}{u}$ ;  $\frac{\Delta r}{r_u} \sim \frac{1}{T_{2u}}$ ),

and  $\Delta l = 1$  (i.e., right symmetry for non-zero matrix element)

$$d = a_0 e \bar{u}^2 / \sqrt{2} = 1776 \text{ e a}_0$$

(for orbit on xy-Plane & light along z:



Selection rule for dipole transitions (from symmetries):

$$\Delta l = \pm 1; \Delta m = 0, \pm 1$$

$\xrightarrow{\text{field in xy}}$  field in z

Prob1: Many degenerate levels for given  $n \rightarrow$  small perturbations induce transitions  $|u, l=u-1\rangle \rightarrow |u, l=u-2\rangle$  etc!

Soln2: Apply electric field along z:

$\Rightarrow$  energy splitting of degen. states: (linear Stark effect)

$\Rightarrow |uC\rangle = |u, l=u-1, m=0\rangle$  non-degenerate w/  $|u, l=u-2, m=0\rangle$   
 $\Rightarrow$  protection against noise (together w/ selection rules)

Quadratic Stark effect: Ext. field E polarizes atom,  $D = p \cdot E \Rightarrow$

$\Rightarrow$  energy shift  $\propto p E^2$ . For  $|uC\rangle$ , p dep. on  $u$

$\Rightarrow$  can be used to define transition  $|uC\rangle \leftrightarrow |u+1, C\rangle$   
by applying z field.

## Stability of Rydberg States:

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Only possible decay:  $|57C\rangle \rightarrow |50C\rangle$  ek. (selection rules)

Quasi-class. theory: accel. charge radiates, decay time  $\propto$  time to cent  $E_{57} - E_{50}$ . (Note: circular orbits are least accel.  $\Rightarrow$  most stable!)

Result:  $\Gamma_a = \frac{2}{3} \frac{\omega^3}{n^2} \omega_{\text{eg}}$ ;  $\Gamma_n = \frac{1}{T_n}$ , nuclear lifetime

$$\Gamma_a = 28 \text{ s}^{-1} \Rightarrow T_a = 36 \text{ ms}$$

"quality factor"  $Q = \frac{1/T}{1/\omega_{\text{eg}}} \approx 10^{10}$  ( $\approx$  # of oscillations before decay)

Note: Thermal background radiation reduces lifetime by factor  $(1 + u_{\text{th}})$  [ $u_{\text{th}} = \#$  thermal photons, cf. Rabi freq.  $\Omega_a = \Omega_0 \sqrt{1+u}$ !]

$\Rightarrow$  shielding/cooling important!

Which  $n$  should we choose?

Large  $n$ :

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- Larger lifetime
- Larger dipole moment

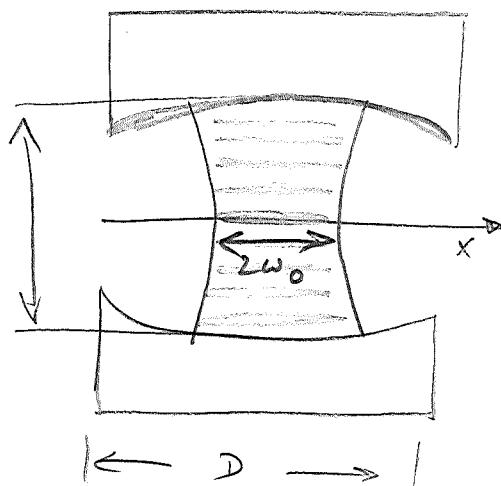
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- polarizability large  $\Rightarrow$  more sensitive to background fields
- Thermal background radiation at  $\omega_{\text{eg}}$  increases.

$\Rightarrow n \approx 51$ .

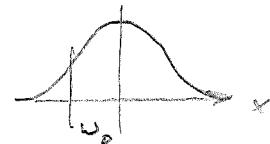
# The cavity: "Fabry-Pérot cavity"

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- \* supercold, Niobium mirrors @ 0.6K
- \* spherical mirror, radius  $r = 40\text{mm}$
- \* diameter  $D = 50\text{ mm}$
- \* distance  $L = 27\text{ mm}$
- \* width  $w_0 = 7.5\text{ mm}$

Mode: TEM<sub>900</sub>: transverse standing wave w/ 9 anti-nodes along  $z$  axis, angular momentum 0, and Gaussian radial intensity profile:



$L$  can be free-fried  $\Rightarrow$  cavity resonant w/ atom,  $\omega_r = \omega_{\text{eg}}$ !

Mode volume:  $V = \frac{\pi}{4} w_0^2 L \approx 700\text{ mm}^3$

$$\Rightarrow E_0 = \sqrt{\frac{tw_0}{2\varepsilon_0 V}} = 1.5 \cdot 10^{-3} \text{ V/m}$$

$\Rightarrow$  macroscopic value of field/photon!

Quality factor of mirrors:  $Q = 3 \cdot 10^8$  (new cavities are  $10^{10}$ )

(i.e.: light is reflected  $3 \cdot 10^8$  times before lost)

$\Rightarrow$  lifetime of photon in cavity  $T = \frac{Q}{\omega_r} \approx 1\text{ ns}$

Avg. photon # in cavity @ 0.6K: 0.06 thermal photons

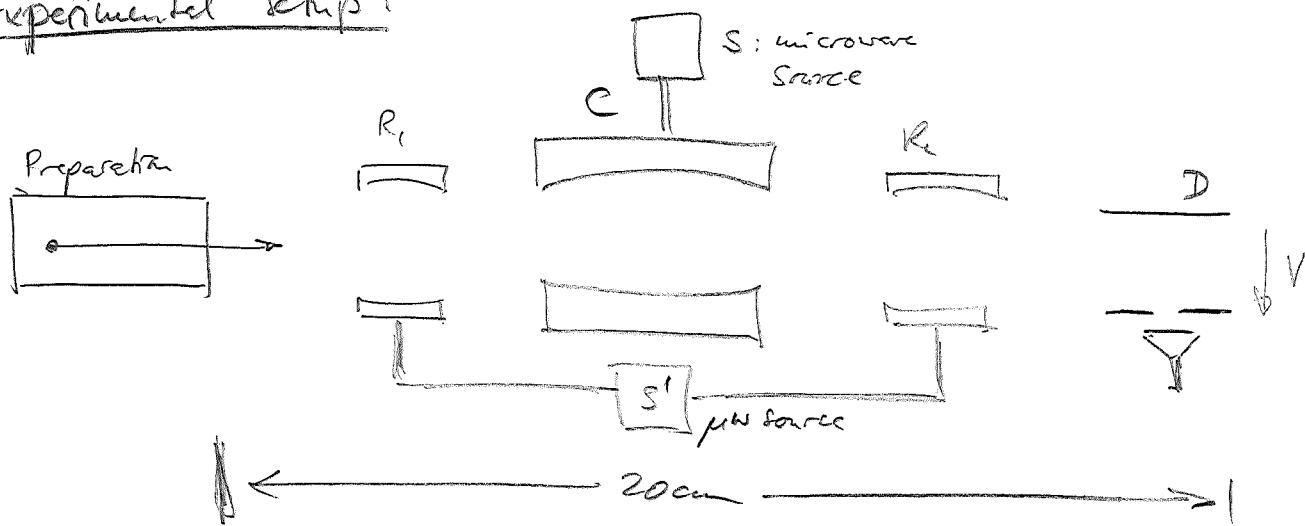
## Atom - Cavity coupling

$$\left. \begin{array}{l} d = 1776 \text{ ea}_0 \\ \epsilon_0 = 1.5 \cdot 10^{-3} \text{ V/m} \end{array} \right\} \Rightarrow \Omega_0 / 2\pi = 50 \text{ kHz} \approx 20 \mu\text{s}$$

### Time scales:

field: $\omega_{\text{cav}}/2\pi$	interaction: $\Omega_0/2\pi$	cavity decay $T_{\text{cav}}$	atom decay $T_a$	in Hz
$5 \cdot 10^{10}$	$5 \cdot 10^4$	$1 \cdot 10^3$	$3 \cdot 10^1$	

### Experimental setup:



Preparation: Prepares pulses w/  $\sim 0.1$  atom/pulse with well-defined speed. (Details: Book by Kerman & Leibundgut)

- Speed  $140-600 \text{ m/s} \pm 2 \text{ m/s}$  (Note:  $2\pi$ -interference in cavity  $\sim 400 \text{ m/s}$ )
- time to pass system  $\sim 1 \text{ ms}$
- spreading of msec. atoms negligible
- Works by selectively exciting atoms w/ specific speed in atom beam

Atoms prepared in  $|e\rangle$  or  $|g\rangle$ .

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R<sub>1</sub>/R<sub>2</sub>: "bad cavity" - applies classical field to atom.

→ Rotates atom via Ramsey pulse,

$$T\text{-pulse} \propto e^{i\pi\delta_x t} \Rightarrow \text{Rotates } |e\rangle \rightarrow |g\rangle \\ |g\rangle \rightarrow -|e\rangle$$

$$\eta_2\text{-pulse} \propto e^{i\eta_2\delta_x t} \Rightarrow \text{Rotates } |e\rangle \rightarrow (|e\rangle + |g\rangle)/\sqrt{2} \\ |g\rangle \rightarrow (|e\rangle - |g\rangle)/\sqrt{2}$$

Used to prepare atom & to rotate for read-out.

C: Cavity. Interaction time tuned by speed or by detuning atom via electric field in 2 directions.

Cavity coupled to (weak) microwave field via small hole in cavity.

D: Detection by ionization in cl. field & amplifying electron.

Trapping field  $\Leftrightarrow$  Rydberg state ( $n=50$ : 195 V/m;  $n=51$ : 134 V/m)

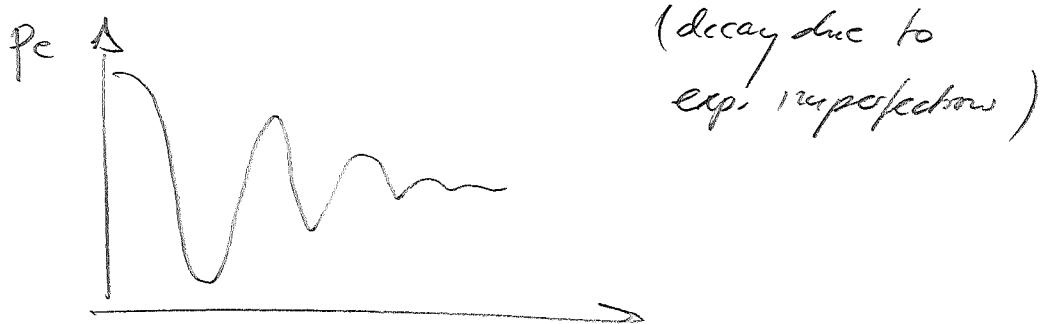
Since # states  $\approx 0.1$ , postselection of results is necessary.

# Experiment 1: Rabi oscillations

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- Atom prepared in  $|e\rangle$ .
- Cavity empty :  $|0\rangle$
- No detuning :  $\omega_c = \omega_{cg}$
- No pulse in  $R_1$
- No pulse in  $R_2$ : atom measured in  $\{|e\rangle, |g\rangle\}$  basis.

$\Rightarrow$  Vacuum Rabi oscillations:



- Cavity initialized in weak coherent field (inj. from microwave source):

