

Variational optimization over MPS:

Find best MPS approx. to G.S. of $H = \sum h_i$.

Use: $|\psi\rangle = \boxed{A^{[1]}} - \boxed{A^{[2]}} - \dots - \boxed{A^{[N]}}$

$\rightarrow |\psi\rangle$ is linear in each $A^{[i]}$

\Rightarrow minimization of $E(A^{[i]}) = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle}$ eigenvalue prob.

Important trick: Isometric gauge:



Today: more on numericals w/ MPS:

① Truncating the bond dimension

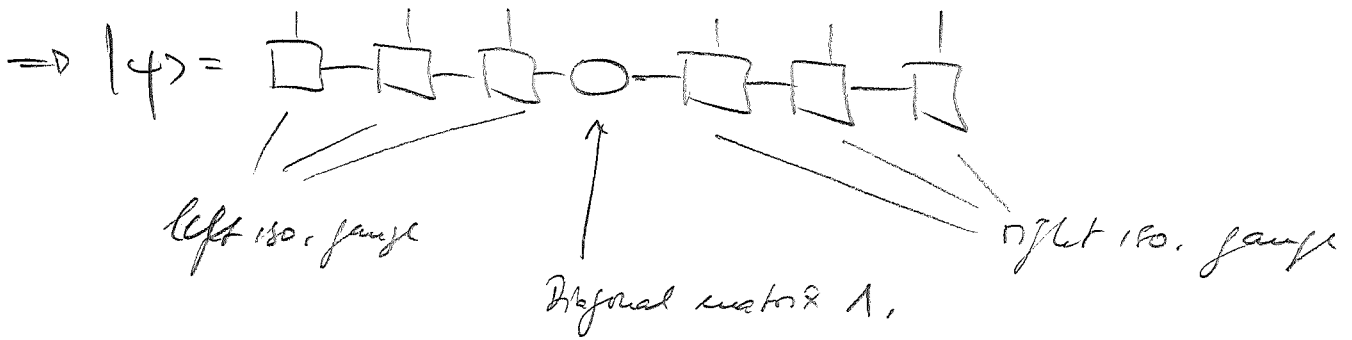
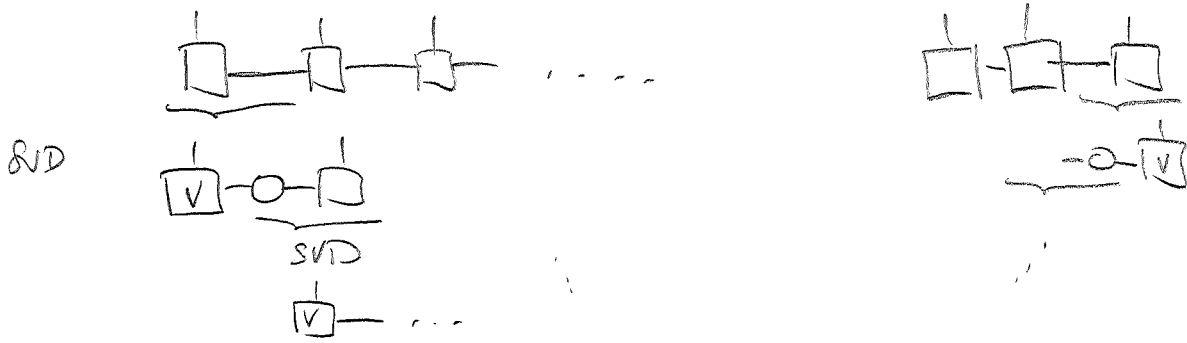
Given MPS $|\psi\rangle$ of some D , can we approximate it by an MPS $|\phi\rangle$ with $D' < D$?

Method I: Do it as a DMRG - optimize $B^{[i]}$ (n/p) sequentially;

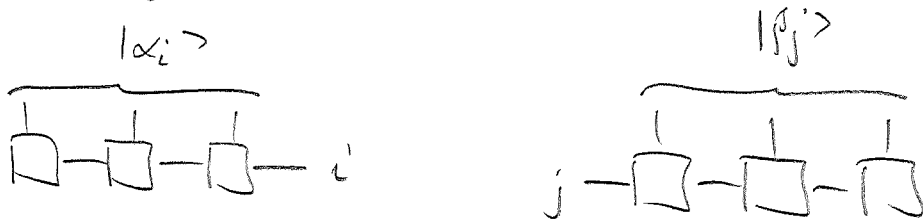
s.t. $\frac{|\langle \psi | \phi \rangle|^2}{\langle \phi | \phi \rangle}$ becomes maximal \Rightarrow greedy, problem!

Method II: Truncation

Choose cut s & choose isometric gauge:



Regard left/right as bases:



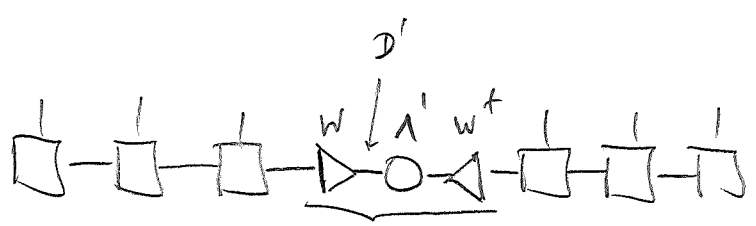
$$\langle \alpha_j | \alpha_i \rangle = \begin{matrix} \square & \square & \square & -j \\ | & | & | \\ \square & \square & \square & -i \end{matrix} \stackrel{\text{iso.}}{=} \begin{bmatrix} j \\ i \end{bmatrix} = \delta_{ij}$$

$\Rightarrow |\alpha_i\rangle, |\beta_j\rangle$ orthonormal bases!

\Rightarrow Schmidt decomposition $|\psi\rangle = \sum_{i=1}^D \lambda_{ii} |\alpha_i\rangle |\beta_i\rangle$

Truncation: Assume wlog $\lambda_{11} \geq \lambda_{22} \geq \dots$

\Rightarrow cut him at $D' < D$.



\triangleleft_W : Isometry $(\begin{matrix} \lambda_{11} & & \\ & \ddots & \\ & & \lambda_{D'D'} \end{matrix} | 0 \rangle)$. $\triangleleft = (\triangleleft_W)^\dagger$

$$\lambda' = \begin{pmatrix} \lambda_{11} & & \\ & \ddots & \\ & & \lambda_{D'D'} \end{pmatrix}$$

\Rightarrow $\triangleleft_W \lambda' \triangleleft \approx \lambda$, and new sub dim. is D' .

What is the error?

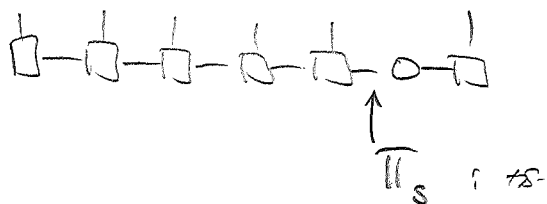
$$\left\| \sum_{i=1}^D \lambda_{ii} |\alpha_i\rangle\langle\beta_i| - \sum_{i=1}^{D'} \lambda_{ii} |\alpha_i\rangle\langle\beta_i| \right\|$$

$$= \left(\sum_{i=D'+1}^D (\lambda_{ii})^2 \right)^{1/2} =: \epsilon_S \quad (\text{"tail weight"})$$

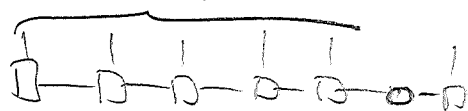
Perform this at all cuts.

→ Maximal error is $\epsilon = \sum \epsilon_s$.

Proof: Consider left no gauge. Cut = Projection on Bond:



no gauge: $\Pi_s \stackrel{\Delta}{=} P_s$



Determine all cuts $\Pi_s \stackrel{\Delta}{=} P_s$ this way:

$$|\psi_{\text{true}}\rangle = P_1 \cdot P_2 \cdot \dots \cdot P_{N-1} |\psi\rangle$$

(since we can sequentially merge $P_{N-1}, P_{N-2}, \dots, P_1$ to $\Pi_{N-1}, \Pi_{N-2}, \dots$)

And: $|\psi_{N-1}\rangle = P_{N-1} |\psi\rangle = |\psi\rangle + \underbrace{|\delta\psi_{N-1}\rangle}_{\|\cdot\| = \epsilon_{N-1}}$

$$\begin{aligned} P_{N-2} P_{N-1} |\psi\rangle &= P_{N-2} |\psi_{N-1}\rangle = \underbrace{P_{N-2} |\psi\rangle}_{=|\psi\rangle + |\delta\psi_{N-2}\rangle} + \underbrace{P_{N-1} |\delta\psi_{N-1}\rangle}_{\|\cdot\| \leq \|\delta\psi_{N-1}\rangle\| = \epsilon_{N-1}} \\ &= |\psi\rangle + \underbrace{|\delta\psi_{N-2}\rangle}_{\|\cdot\| \leq \epsilon_{N-2}} \end{aligned}$$

⇒ $\| P_1 \dots P_{N-1} |\psi\rangle \leq \|\psi\rangle \| \leq \sum \epsilon_s$ □

Note: "foul weight" ϵ_S is related to entanglement

108

entropy:

$$\epsilon_S^2 = \sum_{D'}^D \underbrace{\lambda_{ii}^2}$$

↑ distr. for entropy!

Area law (for $S_d, d \ll 1$) $\Rightarrow \lambda_{ii}^2$ decay quickly \Rightarrow

$\epsilon_S^2 \rightarrow 0$ fast as $D' \rightarrow \infty \Rightarrow$ efficient approximation.

Note: (i) Faster than seq. optimization, but like might be little.

(ii) Truncation can also be parallelized to transp. invariant quantum systems.

Applications: real/imaginary time evolution:

$$H = \sum h_i \cdot |\psi_0\rangle \text{ ground state}$$

$$|\psi_0\rangle \approx \frac{e^{-\beta H}}{\text{norm}} | + \dots + \rangle$$

↑ or some other product state...

requires dep. on properties (gap etc.) of H .

How can we apply $e^{-\beta H}$ efficiently (in a tensor network picture)?

Trotter (L2-Suzuki-Trotter) decomposition:

Wlog: $H = \sum h_{i,i+1}$

$$\Rightarrow H = \underbrace{\left(\sum_{i=0,2,4,\dots} h_{i,i+1} \right)}_{=: h_{\text{even}}} + \underbrace{\left(\sum_{i=1,3,5,\dots} h_{i,i+1} \right)}_{=: h_{\text{odd}}}$$

→ terms within $h_{\text{even}} / h_{\text{odd}}$ mutually commutative!

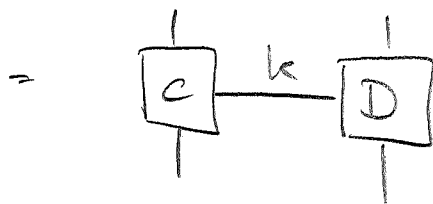
$$\begin{aligned} \Rightarrow e^{-\beta H} &= \left(e^{-\beta H / \pi} \right)^\pi \approx \left(e^{-\beta h_{\text{even}} / \pi} e^{-\beta h_{\text{odd}} / \pi} \right)^\pi \\ &= \prod_{\text{even}} e^{-\beta h_{i,i+1} / \pi} = \prod_{\text{odd}} e^{-\beta h_{i,i+1} / \pi} \end{aligned}$$

→ product of local operators!

→ error like $O\left(\frac{1}{\pi}\right)$ - or better for "high-order" Trotter.

Translate to tensor networks:

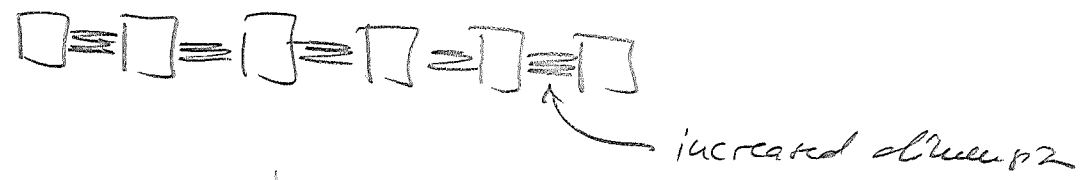
$$\frac{e^{-\beta h_{i,i+1} / \pi}}{\text{two-site operator}} = \sum_k C_k \otimes D_k$$



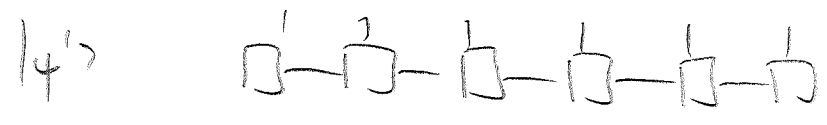
etc.



Block tensors in columns



truncate



→ Iterate until convergence is reached!

Note: This construction also can give an intuition why TSPS approx. G.S. well — it gives a construction for G.S. with large D , but in every ϵ - $\beta_{i,i+1}/M$, the bond is used "very little" (as we are very close to \mathbb{I}) — the bond can be compressed well. (→ can be turned into a proof)

Similar idea for real time evolution:

106

$$|\psi(t)\rangle = e^{-iHt} |\psi(0)\rangle$$

→ can use again Trotter, truncation etc

But: ① less stable: imag. time evol. goes to G.S. ⇒ "self-correcting"

real time evol: unitarity ⇒ errors stay forever.

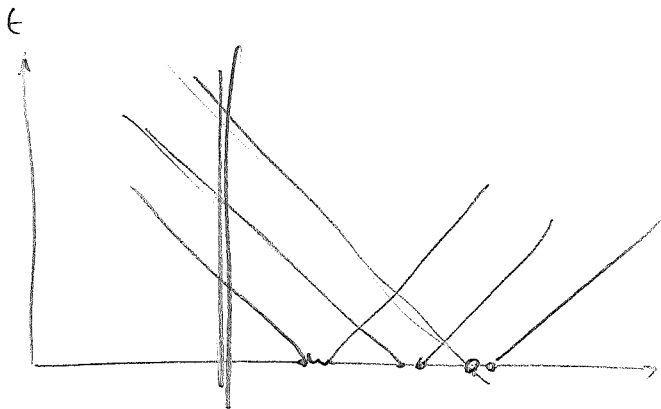
(Test: evolve + un-evolve: $e^{iHt} [e^{-iHt} |\psi(0)\rangle] \stackrel{?}{=} |\psi(0)\rangle$.)

Major bottleneck: In time evolution, entropy typ. grows linearly

with time: $S(t) \sim t$; i.e., $D \sim \exp(t)$.

Intrinsic: Initial state has constant density of excitations.

Excit. are very entangled + move in opposite directions.



≠ excitations through cut $\sim t$!