

(Can be interpreted as "database search": want 104 to find "marked element" x_0 in an unstructured database.)

Assume for now that $x_0: f(x_0) = 1$ is unique.
(Generalization: later / homework)

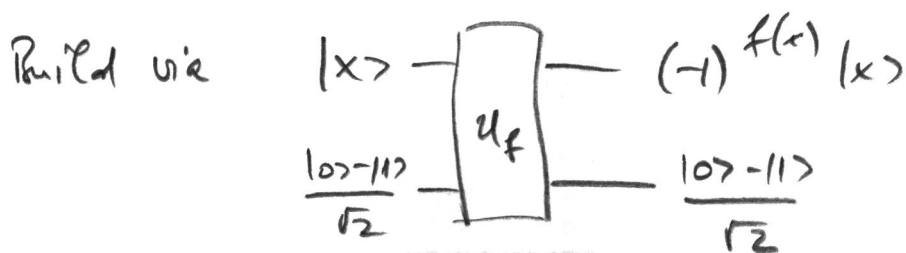
Classically: Need $O(N)$ queries to f for an unstructured search (i.e., w/out using properties of f).

Quantum computers: Will show that $O(\sqrt{N})$ queries enough.

(Note: Only quadratic speedup, but for a very large class of relevant problems)

Ingredient 1:

$$\text{Oracle } O_f : |x\rangle \mapsto (-1)^{f(x)} |x\rangle = (-1)^{\delta_{x,x_0}} |x\rangle$$



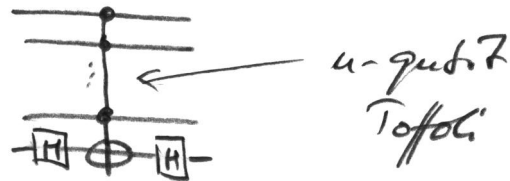
i.e., O_f flips amplitude of "marked" element.

$$\text{Note that } O_f = I - 2 \cdot |x_0\rangle\langle x_0|$$

Ingredient 2:

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Unitary $O_0 : |x\rangle \mapsto (-1)^{\delta_{x,0}} |x\rangle$

Corresponds to $C^{u-1}Z =$ 

\rightarrow can be realized efficiently

Again, $O_0 = I - 2|0\rangle\langle 0|$

Define $O_\omega := H^{\otimes u} O_0 H^{\otimes u} = I - 2|\omega\rangle\langle\omega|$; $|\omega\rangle = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle$.

(Remark: We assumed here $N=2^u$, but not necessary.
Also, every search problem can be triv. embedded
s.t. $N=2^u$.)

Algorithm:

Start from $|\psi_0\rangle = |\omega\rangle = H^{\otimes u} |0\rangle$.

Apply Grover iteration

$$G = -H^{\otimes u} O_0 H^{\otimes u} O_f = -O_\omega O_f;$$

$$|\psi_k\rangle \mapsto |\psi_{k+1}\rangle = G |\psi_k\rangle = -O_\omega O_f |\psi_k\rangle.$$

Observation: Only 2 "special" vectors in O_f, O_w : 106

$|x_0\rangle$ and $|w\rangle \Rightarrow$ can analyze everything in two-dim. space spanned by $|x_0\rangle$ and $|w\rangle$!

Define $|\alpha\rangle := |x_0\rangle$

$$|\beta\rangle := \frac{1}{\sqrt{N-1}} \sum_{x \neq x_0} |x\rangle \propto |w\rangle - \frac{1}{\sqrt{N}} |x_0\rangle \quad \left. \vphantom{\sum} \right\} |\alpha\rangle \perp |\beta\rangle$$

We can always rewrite

$$a|\alpha\rangle + b|\beta\rangle = x|w\rangle + y|w^\perp\rangle, \text{ with } |w^\perp\rangle \perp |w\rangle$$

What is effect of O_f and $(-O_w)$?

$$O_f (a|\alpha\rangle + b|\beta\rangle) \stackrel{\uparrow}{=} -a|\alpha\rangle + b|\beta\rangle$$

$$O_f = I - 2|\alpha\rangle\langle\alpha|$$

\Rightarrow Reflection about $|\beta\rangle$!

$$(-O_w)(x|w\rangle + y|w^\perp\rangle) = x|w\rangle - y|w^\perp\rangle$$

\rightarrow Reflection about $|w\rangle$!

i.e.: Grover iteration = 1) reflect about $|\beta\rangle$

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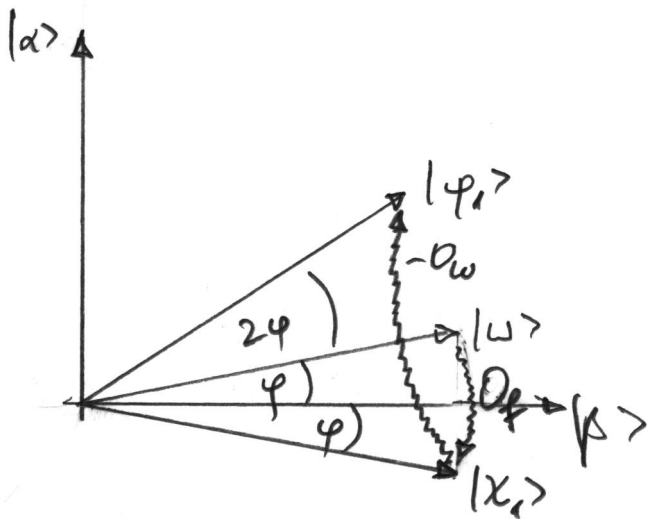
2) reflect about $|\omega\rangle$

So what happens in one iteration, if we start with $|\psi_0\rangle = |\omega\rangle$?

$$|\omega\rangle = \sin\varphi |\alpha\rangle + \cos\varphi |\beta\rangle$$

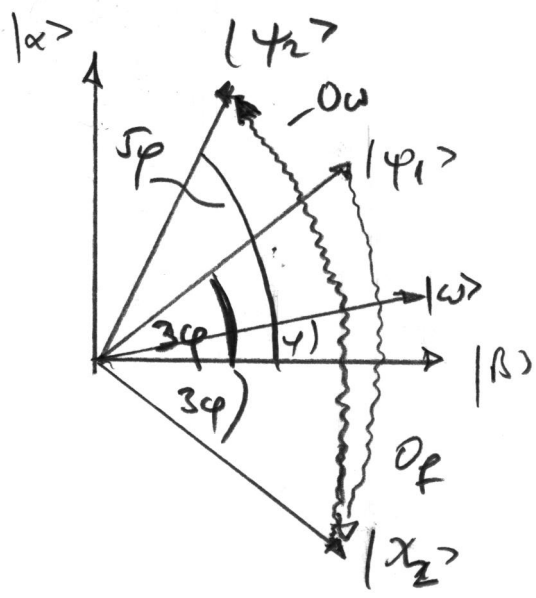
$$|\chi_1\rangle = O_f |\omega\rangle$$

$$|\psi_1\rangle = -O_\omega |\chi_1\rangle = -O_\omega O_f |\omega\rangle$$



$$|\psi_1\rangle = \sin(3\varphi) |\alpha\rangle + \cos(3\varphi) |\beta\rangle.$$

Next iteration: $|\psi_2\rangle = -O_\omega \underbrace{O_f |\psi_1\rangle}_{=: |\chi_2\rangle}$



$$\Rightarrow |\psi_2\rangle = \sin(5\phi) |\alpha\rangle + \cos(5\phi) |\beta\rangle$$

$$\Rightarrow |\psi_k\rangle = \sin((2k+1)\phi) |\alpha\rangle + \cos((2k+1)\phi) |\beta\rangle.$$

Want that $\psi_k = (2k+1)\phi \approx \frac{\pi}{2}$. Then, meas.

will w/ high prob. yield $|\alpha\rangle = |k_0\rangle!$

We have: $|\omega\rangle = \frac{1}{\sqrt{N}} |\alpha\rangle + \sqrt{\frac{N-1}{N}} |\beta\rangle$
 $= \sin\phi |\alpha\rangle + \cos\phi |\beta\rangle$

$$\Rightarrow \frac{\sin\phi}{\cos\phi} = \frac{\sqrt{\frac{1}{N}}}{\sqrt{\frac{N-1}{N}}} = \frac{1}{\sqrt{N-1}}$$

$$\Rightarrow \phi \approx \frac{1}{\sqrt{N}} \text{ for large } N.$$

$$\Rightarrow \text{used } k \approx \frac{\pi}{4} \sqrt{N}$$

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$\Rightarrow O(\sqrt{N})$ calls to f sufficient!

Quadratic speed-up w.r.t. classical algorithms
for several real problems!

Note: • K solutions: Same method works with

$$O\left(\sqrt{\frac{N}{K}}\right) \text{ steps } (\rightarrow HW)$$

• Can be adopted to case where K is unknown.

IV.4. The quantum Fourier transform, period finding, and Shor's factoring algorithm

Recall: Simon's algorithm \rightarrow use $H^{\otimes n} \hat{=}$ Fourier trafo
over $\mathbb{Z}_2^{\otimes n}$ to find period in $\mathbb{Z}_2^{\otimes n}$.

\rightarrow Can we construct a general Q. Fourier Trafo?

\rightarrow Can it be implemented efficiently?

\rightarrow Applications?

a) The Quantum Fourier Transform (QFT)

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Fourier transform (FT) on \mathbb{C}^N :

$$x = (x_0, \dots, x_{N-1}) \in \mathbb{C}^N$$

$$y = (y_0, \dots, y_{N-1}) \in \mathbb{C}^N$$

$$\text{FT: } x \mapsto y \text{ s.t. } y_k = \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} x_j \cdot e^{2\pi i jk/N}$$

Define QFT:

$$|j\rangle \mapsto \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{2\pi i jk/N} |k\rangle$$

$$\text{(Note: QFT: } \sum x_j |j\rangle \mapsto \sum x_j e^{2\pi i jk/N} |k\rangle = \sum y_k |k\rangle$$

\Rightarrow QFT applies FT to amplitudes)

Computational cost of FT: $O(N^2)$ operations.

With $N=2^n \rightarrow$ exp. cost in # of bits n .

Fast FT (FFT): $O(N \log N)$, but still $n \exp(n)$.

Will show: QFT can be implemented in $O(n^2)$ steps

\rightarrow exponential speed-up!

Rewrite QFT in binary:

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* $N = 2^n$

* Write j in binary: $j = j_1 j_2 \dots j_n = j_1 2^{n-1} + j_2 2^{n-2} + \dots + j_n 2^0$

* Decimal part: $0.j_1 j_2 \dots j_n = \frac{1}{2} j_1 2^{-1} + \frac{1}{4} j_2 2^{-2} + \dots$

Then:

$$|j\rangle \mapsto \frac{1}{2^{n/2}} \sum_{k=0}^{2^n-1} e^{2\pi i j k / 2^n} |k\rangle$$

$$= \frac{1}{2^{n/2}} \sum_{k_1=0,1} \dots \sum_{k_n=0,1} e^{2\pi i j \left(\sum_{\ell=1}^n k_\ell 2^{-\ell} \right)} |k_1, \dots, k_n\rangle$$

$$= \bigotimes_{\ell=1}^n \left[\frac{1}{\sqrt{2}} \sum_{k_\ell=0,1} e^{2\pi i j k_\ell 2^{-\ell}} |k_\ell\rangle \right]$$

$$= \bigotimes_{\ell=1}^n \frac{1}{\sqrt{2}} \left[|0\rangle + e^{2\pi i j 2^{-\ell}} |1\rangle \right]$$

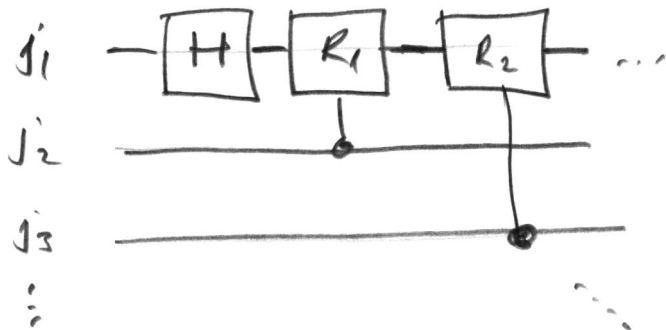
$$\rightarrow j 2^{-\ell} = \underbrace{j_1 j_2 \dots j_{n-\ell} \cdot j_{n-\ell+1} \dots j_n}_{e^{2\pi i \cdot \text{integer}} = 1}$$

$$= \frac{|0\rangle + e^{2\pi i 0 \cdot j_n} |1\rangle}{\sqrt{2}} \otimes \frac{|0\rangle + e^{2\pi i 0 \cdot j_{n-1} j_n} |1\rangle}{\sqrt{2}} \otimes \dots \otimes \frac{|0\rangle + e^{2\pi i 0 \cdot j_1 \dots j_n} |1\rangle}{\sqrt{2}}$$

How to implement this map?

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Start w/ rightmost term: $\frac{|0\rangle + e^{2\pi i \theta_j j_1 j_2 \dots j_u} |1\rangle}{\sqrt{2}}$



$$R_d = \begin{pmatrix} 1 & 0 \\ 0 & e^{2\pi i \cdot 2^{-(d+1)}} \end{pmatrix}$$

action of circuit (up to control):

$$H: |j_1\rangle \mapsto \frac{|0\rangle + e^{2\pi i \theta_j j_1} |1\rangle}{\sqrt{2}}$$

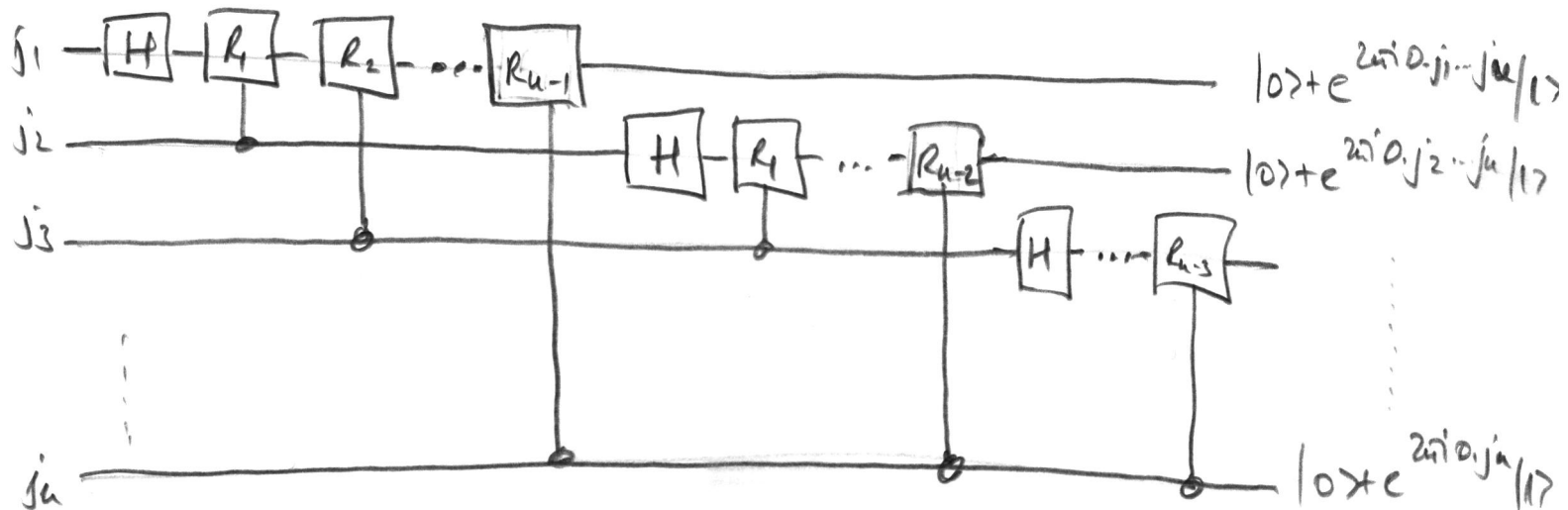
$$C-R_1: \left(\frac{|0\rangle + e^{2\pi i \theta_j j_1} |1\rangle}{\sqrt{2}} \right) |j_2\rangle \mapsto \left(\frac{|0\rangle + e^{2\pi i \theta_j j_1 j_2} |1\rangle}{\sqrt{2}} \right) |j_2\rangle$$

$$C-R_2: \left(\frac{|0\rangle + e^{2\pi i \theta_j j_1 j_2} |1\rangle}{\sqrt{2}} \right) |j_2\rangle |j_3\rangle \mapsto \left(\frac{|0\rangle + e^{2\pi i \theta_j j_1 j_2 j_3} |1\rangle}{\sqrt{2}} \right) |j_2\rangle |j_3\rangle$$

⋮ etc.

⇒ obtain u -th qubit of QFT on 1st qubit.


Continue like that for $(j_2, \dots, j_u), (j_3, \dots, j_u), \dots$



Gate Count: $\frac{n(n+1)}{2} = O(n^2)$ gates!

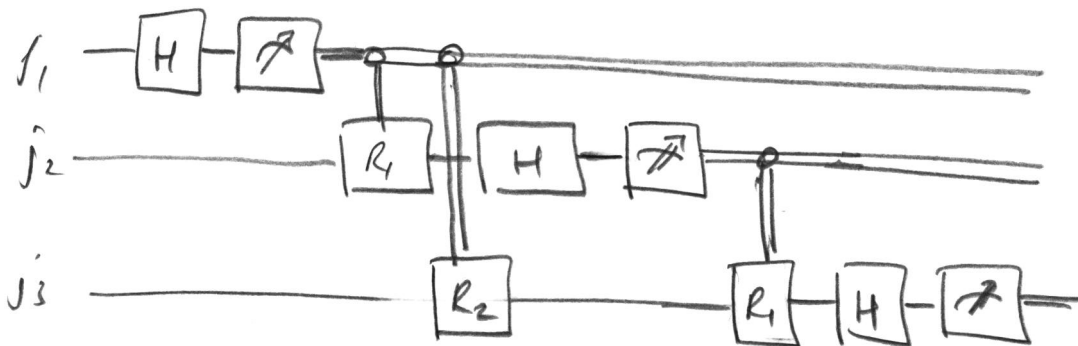
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Notes: • Output is in reverse order (but reordering $\sim O(n)$ ops).

•  \Rightarrow can reverse C-Rd gates.

Upper (control) line acts as control in comp. basis:

If we measure after QFT, we can meas. after H & control Rd-gates classically!



\Rightarrow only one-qubit gates needed!

6) Period finding

Use of QFT: period finding?

Consider $f: \{0, 1\}^n \rightarrow \{0, 1\}^n$, s.t. $\exists r > 0$,

$f(x) = f(x+r)$ (and otherwise $f(x) \neq f(y)$)

Can we find r ?

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Use $U_f: |x\rangle|y\rangle \mapsto |x\rangle|y \oplus f(x)\rangle$

$$\textcircled{1} \frac{1}{2^{u/2}} \sum_A |x\rangle \sum_B |0\rangle \xrightarrow{U_f} \frac{1}{2^{u/2}} \sum_A |x\rangle |f(x)\rangle_B$$

$\textcircled{2}$ Measure $B \rightarrow A$ collapses to

$$\frac{1}{\sqrt{k_0}} \sum_{k=0}^{k_0} |k_0 + kr\rangle$$

(Note: As for Simon, we could omit this step!)

$\textcircled{3}$ Apply QFT & measure in comp. basis:

$$\rightarrow \frac{1}{2^{u/2} \sqrt{k_0}} \sum_k \sum_{l=0}^{2^u-1} e^{2\pi i (x_0 + kr) l / 2^u} |l\rangle$$

$$= \sum_{l=0}^{2^u-1} e^{2\pi i x_0 l / 2^u} \left[\sum_{k=0}^{k_0-1} \frac{1}{2^{u/2} \sqrt{k_0}} e^{2\pi i k r l / 2^u} \right] |l\rangle$$

$=: a_l$

$|a_l|^2$: probability of outcome l .

If $r \ll 2^n$: many values of k & almost periodic (115)

$\Rightarrow |a_e|^2$ peaked around l s.t. $\frac{r^l}{2^n} \approx \text{integer}$

Explicit analysis of a_e shows: w.h.p., we obtain l s.t.

$$\frac{l}{2^n} \approx \frac{s}{r}$$

\hookrightarrow "with high probability"

If $r \ll 2^n$, this can be used to determine $\frac{s}{r}$ w.h.p.

If s, r are coprime (i.e., $\gcd(s, r) = 1$) - this happens with large enough prob., related to density of primes - we can infer r !

\Rightarrow Quantum Algorithm for period finding!

c) Application: factoring

one use of of period finding: factoring.

Given N (not prime) \rightarrow find non-trivial r : $r | N$
 \uparrow
"divides"